



Single Radial/Azimuthal Mode Photonic Crystal Fibers With Anisotropic Elliptical-Hole Lattice Core

メタデータ	<p>言語: English</p> <p>出版者: IEEE</p> <p>公開日: 2017-10-18</p> <p>キーワード (Ja):</p> <p>キーワード (En): Photonic crystal fiber, single radial/azimuthal mode transmission, elliptical-hole lattice, finite element method</p> <p>作成者: ZHENG, Zhong, 辻, 寧英, EGUCHI, Masashi</p> <p>メールアドレス:</p> <p>所属:</p>
URL	http://hdl.handle.net/10258/00009484

Single Radial/Azimuthal Mode Photonic Crystal Fibers With Anisotropic Elliptical-Hole Lattice Core

Zheng Zhong, *Student Member, IEEE*, Yasuhide Tsuji, *Member, IEEE*, and Masashi Eguchi, *Member, IEEE*

Abstract—Novel single radial/azimuthal mode photonic crystal fibers (PCFs) with an anisotropic elliptical-hole lattice core (EHL) are proposed and demonstrated numerically. High birefringence is achieved by introducing elliptical holes into the core, which is considered as an effective way to break the symmetry. With a full vectorial finite element method, we can optimize the fiber parameters and also investigate the wave guiding characteristics in detail. A bandwidth of 210 nm (the wavelength range from 1.45 to 1.66 μm) has been achieved for single radial mode transmission by the PCF with an actinomorphic EHL. Meanwhile, a bandwidth of 100 nm (the wavelength range from 1.50 to 1.60 μm) for single azimuthal mode transmission has been also achieved by the PCF with a multi-annulus EHL. It is anticipated that the proposed PCFs would be of very much convenience in mode multiplexing systems.

Index Terms—Photonic crystal fiber, single radial/azimuthal mode transmission, elliptical-hole lattice, finite element method.

I. INTRODUCTION

Single mode fibers (SMFs) play an important role for mode multiplexing in optical communication systems in recent years. SMFs can be classified into several types by the guidance mechanisms, such as conventional SMFs based on the total internal reflection [1], Bragg fibers based on the effect of Bragg reflection [2]–[4], and photonic crystal fibers (PCFs) based on the concept of photonic crystals [5], [6]. Due to the flexible structures and unique properties of PCFs, a lot of attentions have been paid on their research and development. One of the important parameters of PCFs is the birefringence, and there are two main methods to achieve a high birefringence. One is by introducing elliptical holes or other asymmetric air holes into the core [7]–[10], since elliptical holes are effective to break the symmetry [8]. The other is to change the refractive index contrast between the core and cladding by doping a liquid crystal in the core [12] or utilizing other materials such as metal, polymer, plastic, and so on. Due to the high birefringence, lots of single-polarization (single x - or y -polarization) single-mode (single HE_{11} mode) fibers have been demonstrated earlier. However, a single radial or azimuthal mode fiber has never been investigated by the PCFs so far.

In this letter, to our knowledge, we firstly propose the novel single radial/azimuthal mode PCFs, which have an anisotropic elliptical-hole lattice core (EHL). First, we put forward an analysis to confirm that how much of the birefringence can

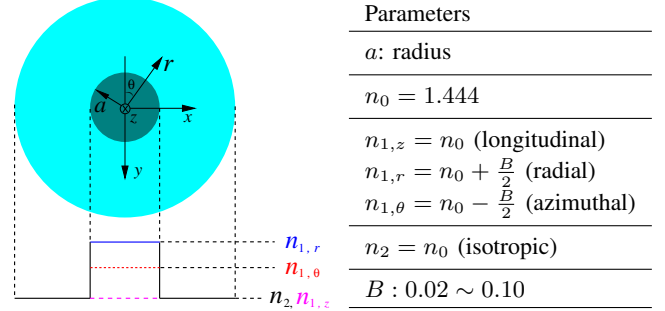


Fig. 1. The schematic and parameters of the equivalent SI-fiber model. The core and cladding are characterized by anisotropic and isotropic materials, respectively. The refractive indices along different directions in the core are assumed in the cylindrical coordinate system. $B(=n_{1,r} - n_{1,\theta})$ represents the birefringence, defined as the refractive index difference between the radial and azimuthal directions of the core.

support the single radial mode transmission based on an equivalent step index- (SI-) fiber model in Section II. Next, we investigate the performances of the single radial mode PCF and single azimuthal mode PCF by employing the full vectorial finite element method (FV-FEM) [13] in Sections III and IV, respectively. Finally, we give a brief conclusion of all analyses in Section V.

II. EQUIVALENT STEP INDEX-FIBER MODEL

For convenience, we consider an equivalent model based on a 2-layer SI-fiber with a radius a of the core, which is shown in Fig. 1. In order to earn a birefringence B , we assume an anisotropic material in the core, whose refractive indices along longitudinal, radial, and azimuthal directions are defined as $n_{1,z} = n_0$, $n_{1,r} = n_0 + \frac{B}{2}$, and $n_{1,\theta} = n_0 - \frac{B}{2}$, respectively. The cladding is characterized by an isotropic material, whose refractive index is defined as $n_2 = n_0$, where n_0 is assumed to be 1.444. At first, we discuss the dispersion of Δn_{eff} as a function of core radius a . Considering a high birefringence of 0.08 obtained in elliptical-hole holey fiber [14], B is changed from 0.02 to 0.10 with an interval of 0.02 and the operating wavelength is set to $\lambda = 1.55 \mu\text{m}$. Δn_{eff} stands for the difference of the modal effective refractive indices between the radial and HE_{11} modes, that can be expressed as the following equation:

$$\Delta n_{\text{eff}} = n_{\text{eff}}^{\text{Radial}} - n_{\text{eff}}^{\text{HE}_{11}}, \quad (1)$$

where $n_{\text{eff}}^{\text{Radial}}$ and $n_{\text{eff}}^{\text{HE}_{11}}$ indicate the modal effective refractive indices of radial and HE_{11} modes, respectively. It can be observed in Fig. 2 that, Δn_{eff} reaches to the maximum of 2.647×10^{-3} under the conditions of $a = 3.5 \mu\text{m}$ and

Z. Zhong and Y. Tsuji are with the Division of Information and Electronic Engineering, Muroran Institute of Technology, Muroran 050-8585, Japan (e-mail: 15043030.y-tsuji@mmm.muroran-it.ac.jp).

M. Eguchi is with the Department of Opto-Electronic System Engineering, Chitose Institute of Science and Technology, Chitose 066-8655, Japan (e-mail: megu@ieee.org).

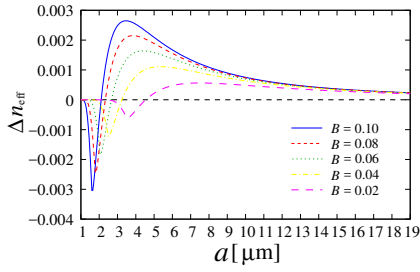


Fig. 2. The difference of the modal effective refractive indices between radial and HE_{11} modes as a function of core radius a , under the operation wavelength $\lambda = 1.55 \mu\text{m}$.

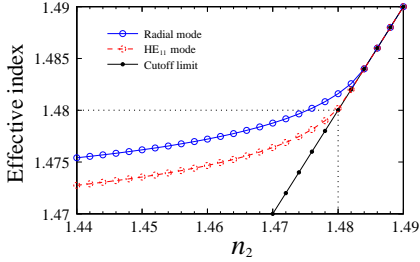


Fig. 3. The modal effective indices of radial and HE_{11} modes with respect to the refractive index of cladding n_2 .

$B = 0.10$. Moreover, Δn_{eff} decreases with an increase in the core radius. In order to achieve the single radial mode transmission, we need to adjust the refractive index of cladding. Fig. 3 shows the modal effective indices of radial and HE_{11} modes with respect to the refractive index of cladding n_2 . When the effective index equals the cladding index, the mode is in cutoff state, and thus we refer to the line with black dots corresponding to the cladding index as cutoff limit. The line with blue-colored open circles and the line with red-colored dashed circles indicate the modal effective indices of radial and HE_{11} modes, respectively. We can observe that HE_{11} mode crosses with the cutoff limit first at $n_2 = 1.480$, and then the radial mode crosses at $n_2 = 1.484$. Therefore, at the refractive index range from 1.480 to 1.484, the single radial mode transmission can be achieved in the equivalent SI-fiber. The magnetic field distributions and Poynting vectors of radial and HE_{11} modes for $n_2 = 1.482$ are shown in Fig. 4. It can be observed that, the HE_{11} mode is not guided and only the radial mode is confined in the core. That means that the single mode transmission can be achieved by the anisotropy between the radial and azimuthal directions of the core in the equivalent SI-fiber. Furthermore, The modal effective indices of radial and HE_{11} modes and the cladding index as a function of wavelength is shown in Fig. 5. The light blue shaded region represents the non-confinement region, where the modal effective indices of radial (blue line) and HE_{11} modes (red dashed line) are lower than that of the cladding index (dot line). In the wavelength range from 1.48 to $1.68 \mu\text{m}$, the single radial mode transmission is achieved in the equivalent SI-fiber.

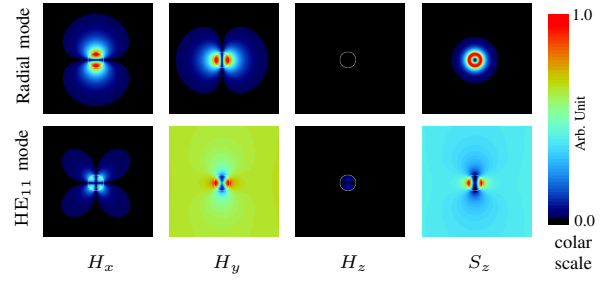


Fig. 4. The magnetic field distributions (H_x , H_y , H_z) and Poynting vectors (S_z) of radial and HE_{11} modes, at $\lambda = 1.55 \mu\text{m}$ and $n_2 = 1.482$.

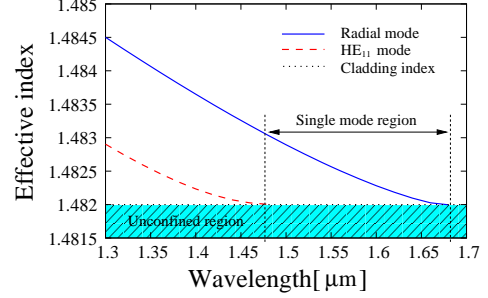


Fig. 5. The modal effective indices of radial and HE_{11} modes and the cladding index as a function of wavelength.

III. SINGLE RADIAL MODE PHOTONIC CRYSTAL FIBER

On the basis of the equivalent SI-fiber model described in Section II, we propose a novel single radial mode PCF with an actinomorphic EHLC to achieve the single radial mode transmission. Fig. 6 shows the schematic of the proposed PCF and the rough mesh division of a quarter cross-section. As shown in Fig. 6 (a), the proposed PCF is characterized by the fused-silica (SiO_2) with defected air holes in a triangular lattice. The centers of all air holes are arranged by the lattice pitch of Λ . The core is arranged by elliptical holes with the radii R_L and R_S for the major and minor axes, respectively. With these major axes of elliptical holes along the radial directions and a large circular hole with a radius of R_a in the center, the radial mode will be guided more easily than the HE_{11} mode. The radius of cladding air hole is represented by R_c . Considering the material dispersion of SiO_2 existing in the real applications, we calculate the refractive index of fused-silica based on the Sellmeier equation [15], which is expressed as follows:

$$n^2(\lambda) = 1 + \frac{0.6961663\lambda^2}{\lambda^2 - (0.0684043)^2} + \frac{0.4079426\lambda^2}{\lambda^2 - (0.1162414)^2} + \frac{0.8974794\lambda^2}{\lambda^2 - (9.8961610)^2}. \quad (2)$$

The refractive index of air is set to 1.0. Due to the structural symmetries along the x - and y -axes, we can only calculate a quarter cross-section to reduce the analysis cost. In Fig. 6 (b), a rough mesh division of a quarter cross-section of the PCF is shown, where the symmetric boundary conditions (Γ_1 , Γ_2) and (Γ_3 and Γ_4) are loaded by the perfect magnetic conductor (PMC) and the perfect electric conductor (PEC), respectively, for the single radial mode PCF analyses. For more accurate

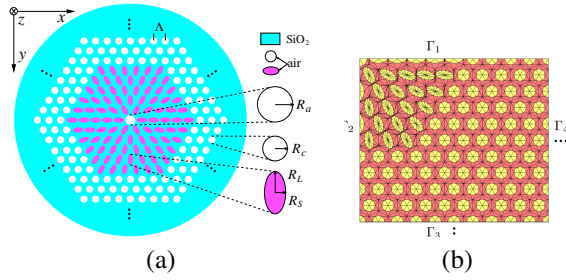


Fig. 6. The schematic of the proposed single radial mode PCF with an actinomorphic EHLC (a), and the rough mesh division of a quarter cross-section (b), where the symmetric boundaries (Γ_1 , Γ_2) and (Γ_3 and Γ_4) are loaded by PMC and PEC, respectively.

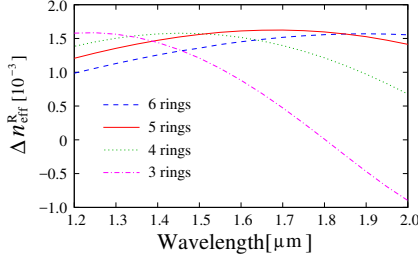


Fig. 7. The modal effective index difference Δn_{eff}^R in different core rings as a function of wavelength.

analyses, we give the rough quarter mesh another 50,000 divisions. After that, we select $\Lambda = 1 \mu\text{m}$, $R_a = 0.40\Lambda$, $R_s = 0.24\Lambda$, and $R_L = 2R_s$ as the optimizing parameters to earn the highest birefringence. First, we set $R_c = \sqrt{2R_s^2}$, which makes both areas of an elliptical hole and an air hole in the cladding equal. In other words, we approximate that the modal effective indices of core and cladding are the same to guide the radial and HE_{11} modes, simultaneously. Fig. 7 shows the relation between the modal effective index difference $\Delta n_{\text{eff}}^R (= n_{\text{eff}}^{\text{Radial}} - n_{\text{eff}}^{\text{HE}_{11}})$ and the number of core rings with respect to wavelength. At the communication wavelength of $1.55 \mu\text{m}$, the maximum of Δn_{eff}^R reaches 1.591×10^{-3} with a 5-ring core. Next, we select an optimal size of R_c to achieve the single radial mode transmission in the proposed PCF. when $R_c = 0.322\Lambda$, the bandwidth of single radial mode transmission is the widest. In the wavelength range from 1.45 to $1.66 \mu\text{m}$ (a bandwidth of 210 nm), the single radial mode transmission is achieved in the proposed PCF with an actinomorphic EHLC. At that time, the magnetic field distributions of both modes are shown in Fig. 8 at a wavelength of $1.55 \mu\text{m}$. The single radial mode is confined, while the HE_{11} mode is not guided. Table I represents the single mode bandwidth (SMB) and its wavelength range (WR) of the proposed PCF, when the elliptical and cladding holes are deviated within $\pm 0.3\%$. when elliptical holes become large (represented by “+”), the single radial mode band shifts toward longer wavelength and the bandwidth is slightly expanded. When they become small (represented by “-”), the band shifts toward to shorter wavelength. On the contrary, the circumstance of cladding holes represents the opposite trend. All the analyses are employed by the FV-FEM with curvilinear hybrid edge/nodal elements, the analyzed region is

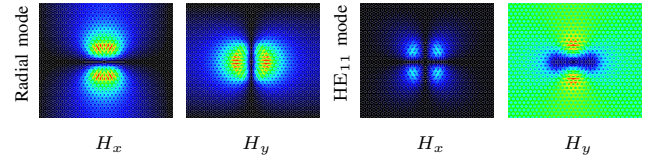


Fig. 8. The magnetic field distributions (H_x and H_y) of both modes in the proposed single radial mode PCF, at $\lambda = 1.55 \mu\text{m}$.

TABLE I
SINGLE RADIAL MODE BANDWIDTH AND ITS WAVELENGTH RANGE VERSUS THE DEVIATION LEVEL OF ELLIPTICAL AND CLADDING HOLES

Deviation level	Elliptical holes		Cladding holes	
	SMB [nm]	WR [μm]	SMB [nm]	WR [μm]
-0.3%	171	1.619~1.790	225	1.262~1.487
-0.2%	182	1.565~1.747	218	1.329~1.547
-0.1%	193	1.510~1.703	212	1.392~1.604
0%	210	1.450~1.660	210	1.450~1.660
+0.1%	209	1.400~1.609	191	1.516~1.707
+0.2%	218	1.340~1.558	175	1.580~1.755
+0.3%	223	1.282~1.505	163	1.638~1.801

about $18.0 \times 15.6 \mu\text{m}^2$, and the total unknowns are 295,550 in the single radial mode PCF analyses.

IV. SINGLE AZIMUTHAL MODE PHOTONIC CRYSTAL FIBER

Based on the discussions in Sections II and III, we consider the schematic of single azimuthal mode PCF with a multi-annulus EHLC as shown in Fig. 9, which is formed by rotating the elliptical holes in the core of Fig. 6 by 90° . As same as the circumstance in the single radial mode PCF analyses, only a quarter cross-section is calculated, and symmetric boundary conditions (Γ'_1 , Γ'_2) and (Γ'_3 , Γ'_4) are loaded by PEC and PMC, respectively, which is shown in Fig. 9 (b). After rotating, the SiO_2 area around the large circular hole gets increased. Fig. 10 shows the modal effective indices of two modes and its difference $\Delta n_{\text{eff}}^A (= n_{\text{eff}}^{\text{Azimuthal}} - n_{\text{eff}}^{\text{HE}_{11}})$ as functions of the radius of large circular hole R_a . Here the radius of cladding air hole is set to $R_c = \sqrt{R_s^2}$. It can be observed that, with the increasing R_a , the effective index of azimuthal mode is almost flat, conversely, the effective index of HE_{11} mode declines. When $R_a = 0.49\Lambda$, Δn_{eff}^A reaches the maximum of 1.205×10^{-3} . On the condition, when we select $R_c = 0.3235\Lambda$, the single azimuthal mode transmission is achieved in the wavelength range from 1.50 to $1.60 \mu\text{m}$, and the bandwidth of that is 100 nm . Fig. 11 shows the magnetic field distributions of both modes at a wavelength of $1.55 \mu\text{m}$. We can observe that the single azimuthal mode has an orthogonal relation with the single radial mode. And the FEM computational window is set to be the same as that in the single radial mode PCF analyses in the previous Section, except that the total unknowns are 295,854 in the single azimuthal mode PCF analyses. Table II represents the single azimuthal mode bandwidth and its wavelength range of the proposed PCF, when the elliptical and cladding holes are deviated only in $\pm 0.1\%$. The phenomenon of the band shift is as same as that in the single radial mode PCF. With a minor change of the air holes, the single azimuthal mode transmission is hardly achieved in the wavelength of $1.55 \mu\text{m}$ in

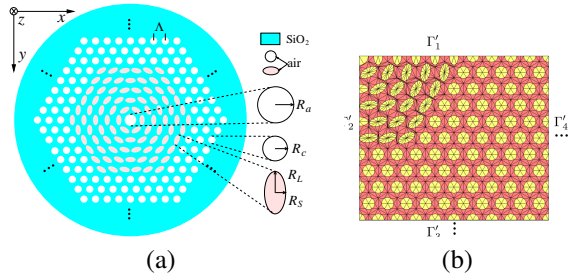


Fig. 9. The schematic of the proposed single azimuthal mode PCF with a multi-annulus EHLC (a), and the rough mesh division of a quarter cross-section (b), where the symmetric boundaries (Γ'_1 , Γ'_2) and (Γ'_3 and Γ'_4) are loaded by PEC and PMC, respectively.

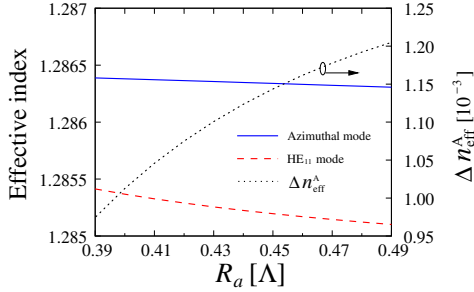


Fig. 10. The modal effective indices of two modes and its difference Δn_{eff}^A as functions of the radius of large circular hole R_a with the radius of cladding air hole $R_c = \sqrt{R_s^2}$.

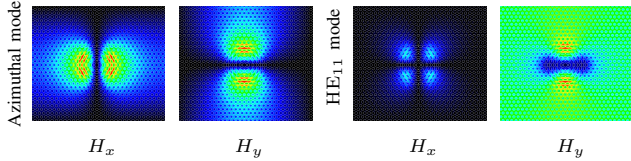


Fig. 11. The magnetic field distributions (H_x and H_y) of both modes in the proposed single azimuthal mode PCF, at $\lambda = 1.55 \mu\text{m}$.

the proposed PCF with a multi-annulus EHLC. The presented results have been estimated for uniform structural fluctuations. So, the presented tolerances are the worst estimations. In fact, the random structural fluctuations may be somewhat canceled out each other for the wavelength range deviations.

V. CONCLUSION

The single radial/azimuthal mode transmission has been numerically demonstrated by the proposed PCFs with an anisotropic EHLC. The analyses based on the equivalent SI-fiber model give the theoretical criteria for designing the proposed PCFs. The bandwidths of 210 nm for the single radial mode transmission and 100 nm for the single azimuthal mode transmission have been achieved in the PCFs with an actinomorphic EHLC and with a multi-annulus EHLC, respectively. Although, the precision problem in fabricating elliptical holes does exist, the research and development of the fabrication techniques are in high progress. Combined with the existing techniques such as the stack and draw technique [16], post-processing technique, flame brushing technique, interlaced stacking technique and so on [17], the fabrication of elliptical holes will not be the most critical problem in the applications. Although the operating wavelength of the

TABLE II
SINGLE AZIMUTHAL MODE BANDWIDTH AND ITS WAVELENGTH RANGE VERSUS THE DEVIATION LEVEL OF ELLIPTICAL AND CLADDING HOLES

Deviation level	Elliptical holes		Cladding holes	
	SMB [nm]	WR [μm]	SMB [nm]	WR [μm]
-0.1%	78	1.563~1.641	106	1.420~1.526
0%	100	1.500~1.600	100	1.500~1.600
0.1%	100	1.430~1.530	75	1.570~1.645

EHLC proposed here is sensitive to structural deviation, we would like to discuss more insensitive structures in the future work. Based on this, we believe that the proposed PCFs will be a significant advance for the mode multiplexing in optical communication systems.

REFERENCES

- [1] T. Okoshi, "Single-polarization single-mode optical fibers," *IEEE J. Quantum Electron.*, vol. QE-17, no. 6, pp. 879–884, Jun. 1981.
- [2] P. Yeh, A. Yariv, and E. Marom, "Theory of Bragg fiber," *J. Opt. Soc. Am.*, vol. 68, no. 9, pp. 1196–1201, Sep. 1978.
- [3] A. Argyros, I. M. Bassett, M. A. V. Eijkelenborg, and M. C. J. Large, "Analysis of ring-structured Bragg fibers for single TE mode guidance," *Opt. Express*, vol. 12, no. 15, pp. 2688–2698, Jun. 2004.
- [4] D. Mao, Z. Ouyang, J. C. Wang, and C. P. Liu, "Single-TM-mode Bragg fibers made of magnetic materials," *Opt. Express*, vol. 16, no. 2, pp. 628–635, Jan. 2008.
- [5] T. A. Birks, J. C. Knight, and P. St. J. Russell, "Endlessly single-mode photonic crystal fiber," *Opt. Lett.*, vol. 22, no. 13, pp. 961–963, Jul. 1997.
- [6] P. St. J. Russell, "Photonic-crystal fibers," *J. Lightw. Technol.*, vol. 24, no. 12, pp. 4729–4749, Dec. 2006.
- [7] K. Saitoh and M. Koshiba, "Single-polarization single-mode photonic crystal fibers," *IEEE Photon. Technol. Lett.*, vol. 15, no. 10, pp. 1384–1386, Oct. 2003.
- [8] M. Eguchi and Y. Tsuji, "Single-mode single-polarization holey fiber using anisotropic fundamental space-filling mode," *Opt. Lett.*, vol. 32, no. 15, pp. 2112–2114, Aug. 2007.
- [9] K. Ichikawa, Z. Zhang, Y. Tsuji, and M. Eguchi, "A single-polarization holey fiber with anisotropic lattice of circular air holes," *J. Lightw. Technol.*, vol. 33, no. 18, pp. 3866–3871, Sep. 2015.
- [10] L. An, Z. Zheng, Z. Li, T. Zhou, and J. Cheng, "Ultrahigh birefringent photonic crystal fiber with ultralow confinement loss using four airholes in the core," *J. Lightw. Technol.*, vol. 27, no. 15, pp. 3175–3180, Aug. 2009.
- [11] S. E. Kim, B. H. Kim, C. G. Lee, S. Lee, K. Oh, and C. S. Kee, "Elliptical defected core photonic crystal fiber with high birefringence and negative flattened dispersion," *Opt. Express*, vol. 20, no. 2, pp. 1385–1391, Jan. 2012.
- [12] M. F. O. Hameed and S. S. A. Obayya, "Modal analysis of a novel soft glass photonic crystal fiber with liquid crystal core," *J. Lightw. Technol.*, vol. 30, no. 1, pp. 96–102, Jan. 2012.
- [13] M. Koshiba and Y. Tsuji, "Curvilinear hybrid edge/nodal elements with triangular shape for guided-wave problems," *J. Lightw. Technol.*, vol. 18, no. 4, pp. 737–743, May 2000.
- [14] M. Eguchi and Y. Tsuji, "Squeezed lattice elliptical-hole holey fiber with extremely high birefringence," *Opt. Lett.*, vol. 33, no. 16, pp. 1792–1794, Aug. 2008.
- [15] N. Karasawa, S. Nakamura, N. Nakagawa, M. Shibata, R. Morita, H. Shigekawa, and M. Yamashita, "Comparison between theory and experiment of nonlinear propagation for a-few-cycle and ultrabroadband optical pulses in a fused-silica fiber," *IEEE J. Quantum Electron.*, vol. 37, no. 3, pp. 398–404, Mar. 2001.
- [16] D. Pysz, I. Kujawa, R. Stepień, M. Klimczak, A. Filipkowski, M. Franczyk, L. Kociszewski, J. Buzniak, K. Harasny, and R. Buczynski, "Stack and draw fabrication of soft glass microstructured fiber optics," *Bull. Pol. Acad. Sci. Tech. Sci.*, vol. 62, no. 4, pp. 667–682, Dec. 2014.
- [17] E. C. Mägi, P. Steinvurzel, and B. J. Eggleton, "Tapered photonic crystal fibers," *Opt. Exp.*, vol. 12, no. 5, pp. 776–784, Mar. 2004.