

Harmonic Resonance on Unbalanced Transmission Lines with a Salient-pole Synchronous Generator

メタデータ 言語: eng

出版者: 室蘭工業大學

公開日: 2014-05-19

キーワード (Ja):

キーワード (En):

作成者: 三浦, 五郎

メールアドレス:

所属:

URL http://hdl.handle.net/10258/3002

Harmonic Resonance on Unbalanced Transmission Lines with a Salient-pole Synchronous Generator

Goro Miura*

Abstract

On account of the saliency of poles of an alternating current generator, a lot of harmonics can be emerged on transmission lines when some kinds of unbalanced faults occur. This paper deals with the mathematical development of a fundamental equation with the case of one line-fault, and clarifies critical conditions of the occurrence of this harmonic resonance.

The author has lately a chance to assist and cooperate Mr. Ogushi, Professor in Hokkaido University, with his research for the problem of harmonic resonance of a transmission line caused by an unbalanced line-fault. As a part of the theoretical development of this research was already published (1), the treatment of mathematics is one of considerably complicated and it would be liable to obstruct the clear understanding of the physical nature of the phenomena.

Soon after, the author succeeded to manage it in the more systematic and more direct procedure by using "tensor"—a powerful mathematical tool in the engineering—and obtained the same results as before; of which process of the treatise is felt valuable and so published here.

As all the winding-axes of a balanced three-phase salient-pole synchronous generator will revolve with synchronous speed because of these holonomic natures, its impedance tensor is represented by a single equation, [Z]=[R]+p[L], according to the Maxwell's equation. If it is assumed [R]=0 as a justifiable approximation,

		f	$oldsymbol{d}_a$	d_b	d_c	q_c	q_b	q_a
$[Z]^{\stackrel{(2)}{=}}p$	f	x_f	x_{af}	x_{af}	xaf	0	0	0
	d_a	x_{af}	x_d	x_{rd}	x_{rd}	0	0	0
	d_b	x_{af}	x_{rd}	x_d	x_{rd}	0	0	0
	d_c	x_{af}	x_{rd}	x_{rd}	x_d	0	0	0
	q_c	0	0	0	0	x_q	x_{rq}	x_{rq}
	q_b	0	0	0	0	x_{rq}	x_q	x_{rq}
	q_a	0	0	0	0	x_{rq}	x_{rq}	x_q

^{*} 三 浦 五 郎
(1). G. Miura: A Prompt Memoir of The Muroran College of Technology Vol. I, No. 2.
(2). The nomenclature owes to: G. Miura; An Analytic Method of the Synchronous Machines, Memoirs of The Muroran College of Technology, Vol. 1, No. 1.

332 G. Miura

On	the	other	hand.	the	current	transformation	tensor	is
-								

		f	a	b	c
	f	1			
	d_a		$\cos \theta_1$		
	d_b			$\cos \theta_2$	
[C]=	d_c			TO A STATE OF THE PARTY OF THE	$\cos \theta_3$
	q_c				$\sin \theta_3$
	q_b			$\sin \theta_2$	
	q_{it}		$\sin \theta_1$		-

Then, a new impedance tensor is calculated from

$$[Z']=[C],[Z][C]=p[L']$$

In the tensor [Z'], the following simplification shall be made:

$$A = \frac{x_d + x_q}{2} , \qquad B = \frac{x_d - x_q}{2} .$$

Then, the above machine constants shall be replaced by the three-phase constants; namely x_{rd} by $\frac{2}{3}(x_d-x_o)$, x_d by $\frac{2}{3}(x_d+\frac{x_o}{2})$,

$$x_{rq}$$
 by $\frac{2}{3}(x_q-x_o)$, x_q by $\frac{2}{3}(x_q+\frac{x_o}{2})$.

If the unnecessary term f is eliminated (assuming the excitation voltage $E_t=0$) by a short-circuit matrix, the final impedance theorem

 $x_a(p)$ and $x_q(p)$ in this case should be equal to x_a' and x_q respectively from the approximation [R]=0. Now, since the steady state phenomena will be discussed here, p=j should be permitted in all above equations (using the per-unit-method, $\omega=1$).

Nextly, it is assumed that the transmission line which is treated here has no resistance, inductance and leakance but has only capacitance between line-to-earth and line-to-line; of which capacitance can be transformed to the equivalent balanced star circuit capacities, one side of capacity being C.

So,
$$\frac{1}{vwC} = -j\frac{1}{C} = -jx_c$$
 is taken.

Then, the phenomena with a line-earth fault can be represented by a next equivalent circuit, Fig. 1.

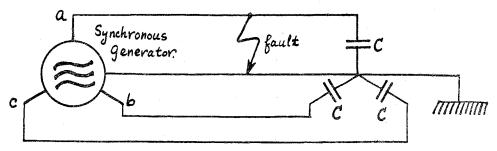


Fig. 1. Equivalent circuit of a line-fault.

The circuit, Fig. 1. may be considered as equal to a circuit Fig. 2,

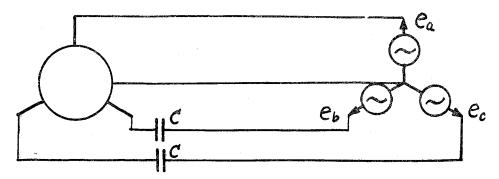


Fig. 2. Equivalent circuit of Fig. 1.

where $e_a = e \sin \theta_1$, $e_b = e \sin \theta_2$, $e_c = e \sin \theta_3$, and exitation D.C. source in a generator is assumed to be absence.

Accordingly, the differential equation will be

$$-e = j(Z - x_o) i.$$

Since j means differentiation, the both sides of the equation should be integrated, resulting

$$\begin{bmatrix} 3 \ e \cos \theta_1 \\ 3 \ e \cos \theta_2 \\ 3 \ e \cos \theta_3 \end{bmatrix} = \begin{bmatrix} x_o + 2A + 2B \cos 2\theta, \\ x_o - A + 2B \cos \overline{2\theta + 120}, \\ x_o - A + 2B \cos \overline{2\theta + 120}, \\ x_o - A + 2B \cos \overline{2\theta + 120}, \\ x_o + 2A + 2B \cos \overline{2\theta - 120} - 3x_c, & x_o - A + 2B \cos 2\theta \\ x_o - A + 2B \cos 2\theta, & x_o + 2A + 2B \cos 2\theta + 120 - 3x_c \end{bmatrix} [i].$$

(21)

334 G. Miura

The inverse calculation may be made as follows:

$$[i] = \frac{e}{D} \left[\begin{array}{ccc} a & b & c \\ g & h & k \\ l & m & n \end{array} \right]$$

where

$$D = 2x_{c} (x_{c} - x_{o}) B \cos 2\theta + (A^{2} - B^{2}) (3x_{o} - 2x_{c}) + x_{c}^{2} (x_{o} + 2A) - 4x_{c}x_{o}A.$$

$$a = A^{2} - B^{2} + 2x_{o}A + 3x_{c}^{2} - 2x_{c} (x_{o} + 2A) + 2B (x_{c} - x_{o}) \cos 2\theta.$$

$$b = A^{2} - B^{2} - x_{o}A - x_{c} (A - x_{o}) + 2B (x_{c} - x_{o}) \cos \overline{2\theta + 120}.$$

$$c = A^{2} - B^{2} - x_{o}A - x_{c} (A - x_{o}) + 2B (x_{c} - x_{o}) \cos \overline{2\theta - 120}.$$

$$h = A^{2} - B^{2} + 2x_{o}A - x_{c} (x_{o} + 2A) - 2B (x_{c} \cos 2\theta + x_{o} \cos \overline{2\theta - 120}).$$

$$k = A^{2} - B^{2} - x_{o}A - 2x_{o} B \cos 2\theta.$$

$$n = A^{2} - B^{2} + 2x_{o}A - x_{c} (x_{o} + 2A) - 2B (x_{c} \cos 2\theta + x_{o} \cos \overline{2\theta + 120}).$$

$$g = b, \qquad l = c, \qquad m = k.$$

Accordingly,

$$-i_a=rac{3e}{D}\left(x_c-x_o
ight)\left(A-B-x_o
ight)\cos heta, \ -i_b=rac{e}{D}\left\{-3x_o\left(A-B
ight)\cos heta_2+x_o\left(A-B
ight)\left(\cos heta_2-\cos heta_3
ight)-x_ox_o\left(\cos heta_1-\cos heta_2
ight)
ight\}. \ -i_c=rac{e}{D}\left\{-3x_o\left(A-B
ight)\cos heta_3+x_o\left(A-B
ight)\left(\cos heta_3-\cos heta_2
ight)-x_ox_o\left(\cos heta_1-\cos heta_3
ight)
ight\}.$$

From the equation of i_a , the resonance condition can be obtained. Namely, if

$$lpha = (A^2 - B^2) (3x_o - 2x_c) + x_c^2 (x_o + 2A) - 4x_c x_o A,$$
 $\beta = -x_c (x_c - x_o) B,$
 $\gamma = 3e (x_c - x_o) (A - B - x_c),$

then

$$i_a = rac{-\gamma \cos \theta}{\alpha - 2\beta \cos 2\theta} = rac{-2\gamma}{\alpha - 2\beta + \sqrt{\alpha^2 - 4\beta^2}} \sum_{n=1, 3, ...}^{\infty} \eta^{\frac{n-1}{2}} \cos n\theta.$$

$$\eta = \frac{1}{2\beta} \left\{ \alpha - \sqrt{\alpha^2 - 4\beta^2} \right\} < 1.$$

When $\alpha-2\beta=0$, i_a diverts to ∞ . The resonance condition is then:

$$(A^2-B^2)(3x_o-2x_c)+x_c^2(x_c+2A)-4x_cx_oA+2x_c(x_c-x_o)B=0.$$

Introducing

$$A = \frac{x_d' + x_q}{2} \quad \text{and} \quad B = \frac{x_d' + x_q}{2},$$

$$x_{c} = \frac{2x'_{d} x_{q} + x_{o} (3x'_{d} + x_{q}) \pm \sqrt{\{2x'_{d} x_{q} + x_{o} (3x'_{d} + x_{q})\}^{2} - 12x'_{d} x_{q} x_{o} (2x'_{d} + x_{o})}}{2(2x'_{d} + x_{o})}.$$

Assuming $x_o \ll x_d'$, $x_o \ll x_q$, the expression inside the root can be approximated. And the final result will be

$$x_c = x_q$$
 or $\frac{3x_o x_d'}{2x_d' + x_o}$.

However, as $x_c = x_q$ will made also, γ , the numerator of i_a , equal to zero, it should be excluded. Accordingly,

$$x_c \bigg/ \frac{3x_o x_d'}{2x_d' + x_o} = n^2$$

where n is the dimension of the harmonics considered. This coincides with the result in the former publication as stated before.

As already stated in the former publication, this research had been made during two months of the last summer vacation at Hokkaido University. The author wishes to express his gratitude for much valuable guidance of Prof. Ogushi of the university, and for much assistance recieved from Dr. Iguchi, the president, and proffesors of Electrical Department, of the Muroran University of Engineering.

(Received March 9, 1950)